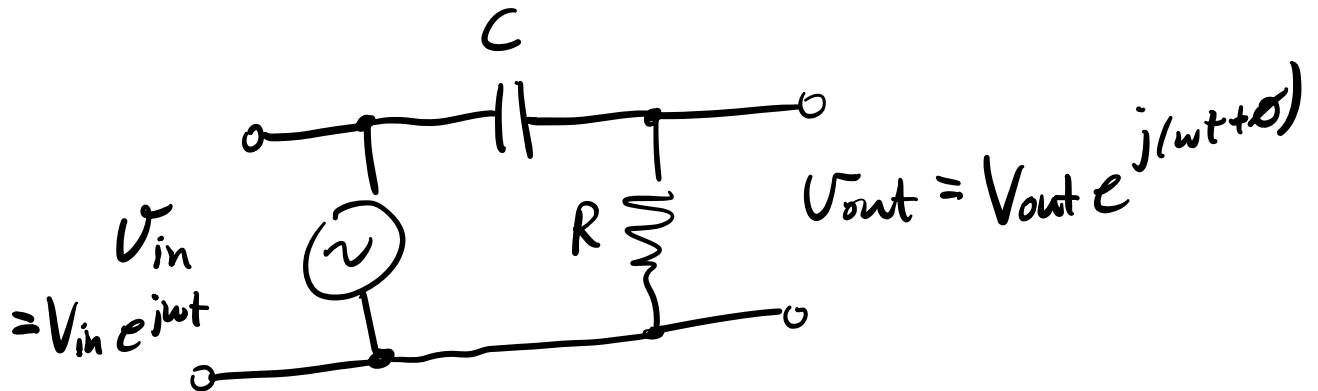


PHYS 331 - Sept. 18, 2023

Last Time:



When V_{out} is out of phase w.r.t. V_{in} ,
then V_{out} has real & imaginary components:

$$X = V_{out} \cos \phi \quad \leftarrow \text{real component}$$

$$Y = V_{out} \sin \phi \quad \leftarrow \text{imaginary "}$$

$$\tan \phi = \frac{Y}{X}$$

$$V_{out} = \sqrt{X^2 + Y^2}$$

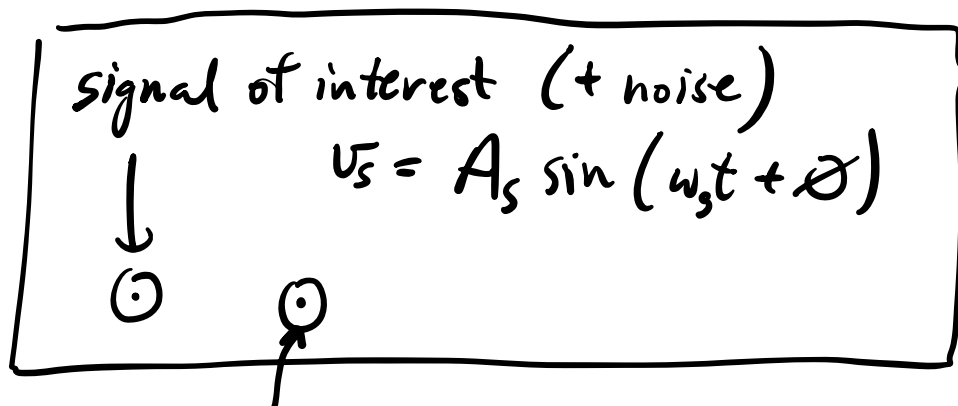
A lock-in detector is able to meas. the real & imag. components of small AC signals even when buried in noise.

If use lock-in to meas. X & Y , we can then find the amp. & phase of the signal

$$\theta = \tan^{-1}\left(\frac{Y}{X}\right) \quad V_{out} = \sqrt{X^2 + Y^2}$$

How does a lock-in detector meas. X & Y ?

Lock-in detectors have two inputs



reference

$$V_r = A_r \sin(\omega_r t + \theta)$$

We know everything about reference v_r .

We control A_r, ω_r, θ

Want to determine A_s & ϕ of v_s .

① Internally, lock-in detector multiplies the two inputs:

$$v_{out} = v_s v_r = A_s A_r \sin(\omega_s t + \phi) \sin(\omega_r t + \theta)$$

$$= \frac{A_s A_r}{2} \left\{ \cos [(\omega_s t + \phi) - (\omega_r t + \theta)] \right.$$

$$\left. - \cos [(\omega_s t + \phi) + (\omega_r t + \theta)] \right\}$$

$$v_{out} = \frac{A_s A_r}{2} \left\{ \cos [(\omega_s - \omega_r)t + (\phi - \theta)] \right.$$

$$\left. - \cos [(\omega_s + \omega_r)t + (\phi + \theta)] \right\}$$

② Next, the multiplied signal is passed through a low-pass filter to eliminate the high-freq. term.

$$V_{out} \rightarrow V_{out}' = \frac{A_s A_r}{2} \cos \left[(\omega_s - \omega_r) t + (\phi - \theta) \right]$$

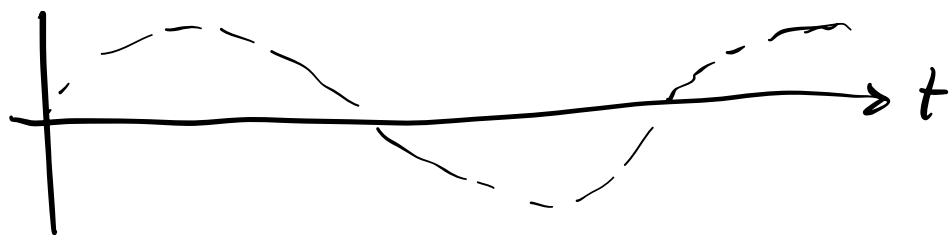
(after filter).

③ The filtered signal is then averaged over a long time.

If $\omega_s \neq \omega_r$, then $\langle V_{out}' \rangle = 0$

time average

b/c $\cos(\omega t)$ averages to zero (half pos. & half neg.)



Lock-in detector is blind to signal w/ freq. $\omega_s \neq \omega_r$.

If $\omega_s = \omega_r$, then

$$\langle v_{out}' \rangle_{\omega_s = \omega_r} = \frac{A_s A_r}{2} \cos(\vartheta - \theta) \quad \text{ⓧ}$$

In ⓧ, we still have two unknowns (A_s, ϑ) and only one eq'n.

④ To establish two eq'ns & two unknowns,

(a) set $\theta = 0$. Set ref. phase arbitrarily to zero.

$$\langle v_{out}' \rangle_{\omega_s = \omega_r} \Big|_{\theta = 0} = \frac{A_s A_r}{2} \cos \vartheta$$

$$= \frac{A_r}{2} \left[\underbrace{A_s \cos \vartheta}_{\text{X}}$$

real component
of v_s

$$X = \frac{2}{A_r} \langle V_{out}' \rangle_{\omega_s = \omega_r} \Big|_{\theta = 0}$$

real component

(b) set $\theta = \frac{\pi}{2}$

$$\langle V_{out}' \rangle_{\omega_s = \omega_r} \Big|_{\theta = \frac{\pi}{2}} = \frac{A_s A_r}{2} \underbrace{\cos\left(\phi - \frac{\pi}{2}\right)}_{\sin \phi}$$

$$= \frac{A_r}{2} \underbrace{\left[A_s \sin \phi \right]}$$

Y imaginary component of v_s

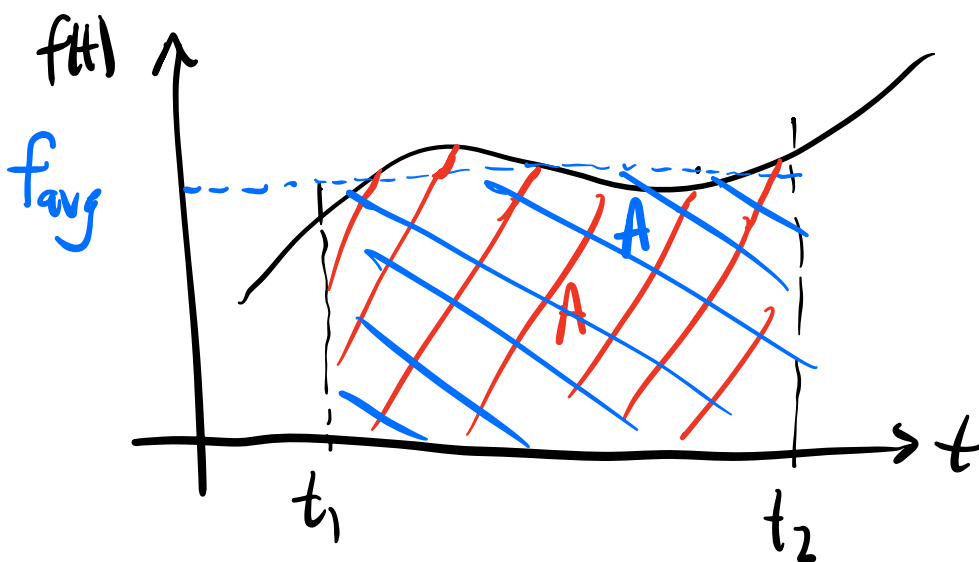
$$Y = \frac{2}{A_r} \langle V_{out}' \rangle_{\omega_s = \omega_r} \Big|_{\theta = \frac{\pi}{2}}$$

Lock-in detector is capable of meas.
 X & Y of our signal of interest V_s .
 Can use these to then determine amp. & phase

$$A_s = \sqrt{X^2 + Y^2} \quad \theta = \tan^{-1} \left(\frac{Y}{X} \right)$$

Suppose that $\omega_s \approx \omega_r$ (but not identical)
 How well does $\langle V_{out} \rangle$ discriminate
 between freq. not equal to ω_r ?

Aside



$$f_{avg} = \frac{1}{t_2 - t_1} \int_{t_1}^{t_2} f(t) dt$$

set red
& blue
areas equal

The average of V_{out}' from $t=0$ to final time T .

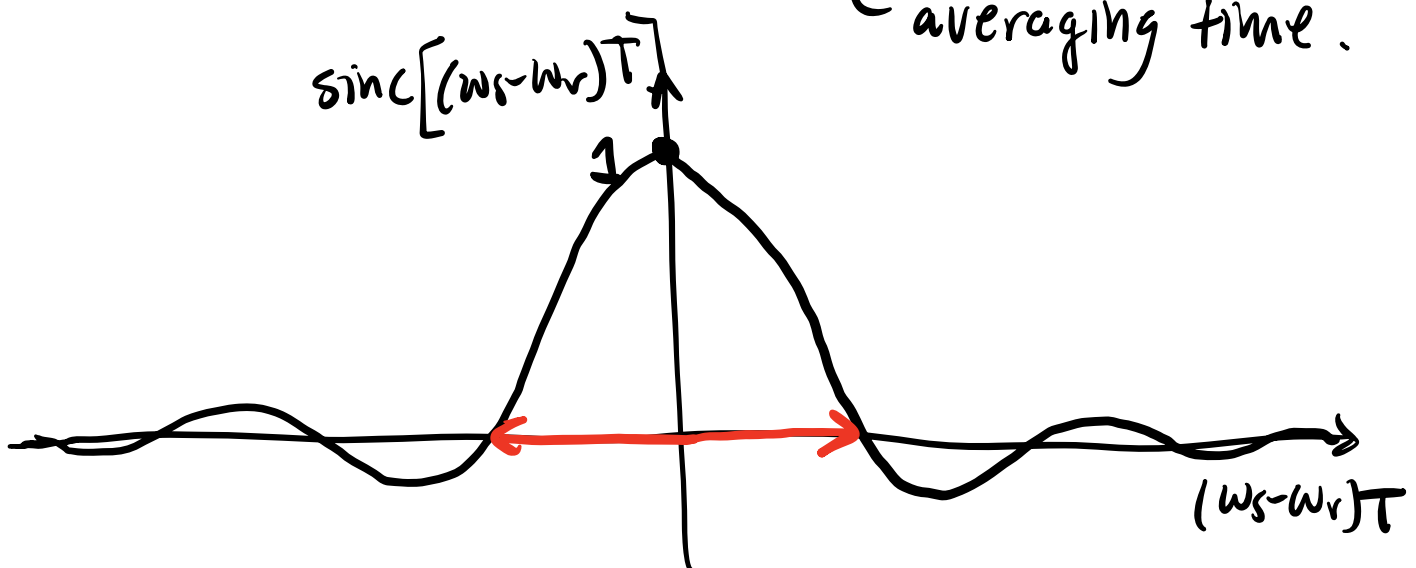
$$\langle V_{out}' \rangle \propto \frac{1}{T} \int_0^T \cos[(\omega_s - \omega_r)t] dt$$

$$\propto \frac{1}{T} \frac{1}{\omega_s - \omega_r} \sin[(\omega_s - \omega_r)t] \Big|_0^T$$

$$\propto \frac{1}{(\omega_s - \omega_r)T} \sin[(\omega_s - \omega_r)T]$$

$$\langle V_{out}' \rangle \propto \text{sinc}[(\omega_s - \omega_r)T]$$

↑
averaging time.



In order to characterize the width of $\langle V_{out} \rangle$, find the values of $(\omega_s - \omega_r)T$ that produce the first zero crossings of $\text{sinc}[(\omega_s - \omega_r)T]$.

The first zeros occur when
 $(\omega_s - \omega_r)T = \pm \pi$

$$T_{\pm} = \frac{\pm \pi}{2\pi \underbrace{(f_s - f_r)}_{\Delta f}}$$

$$\Delta T = T_+ - T_-$$

$$\Delta T = \frac{\pi}{2\pi \Delta f} - \frac{-\pi}{2\pi \Delta f} = \frac{1}{\Delta f}$$

$$\Delta T = \frac{1}{\Delta f}$$

To suppress signal leakage from average, need averaging time to be much larger than $\frac{1}{\Delta f}$.

In our example of an RC circuit, we supplied a sinusoidal input & detected a sinusoidal output.

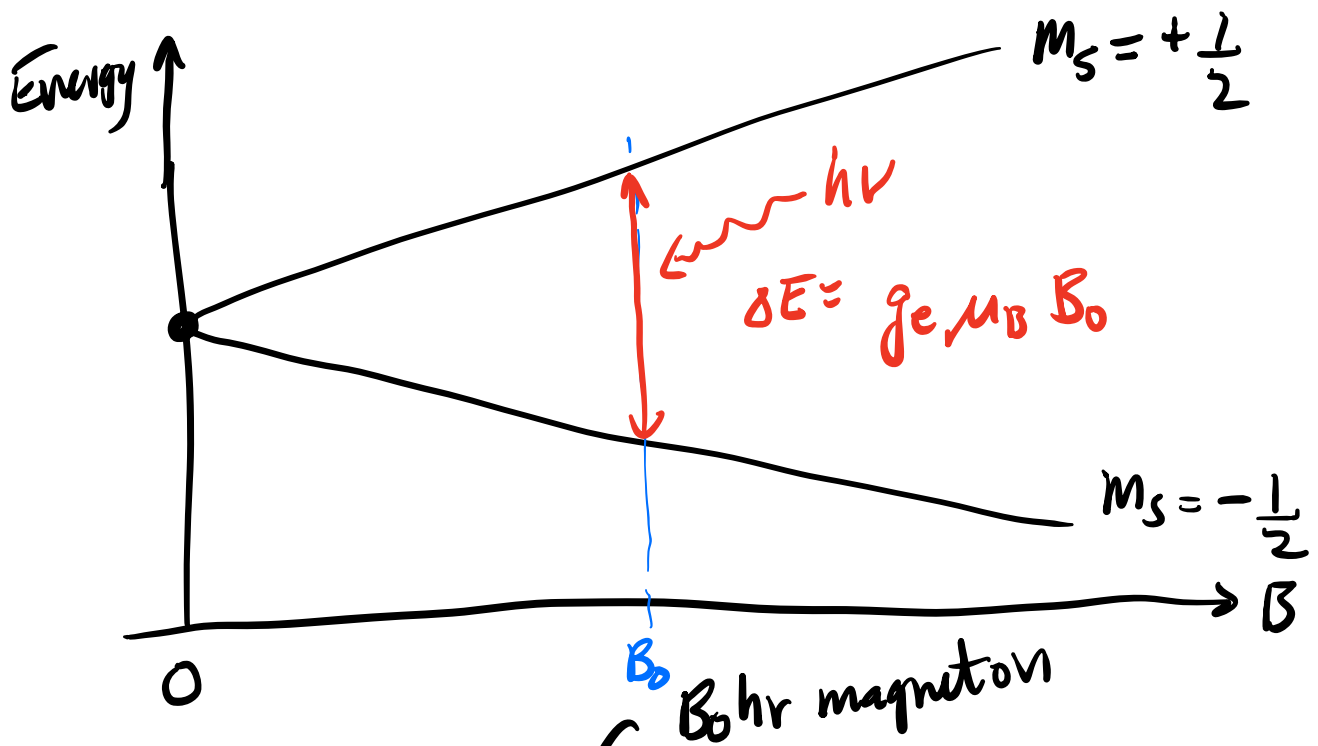
Not all experiments can be driven by a pure sine wave. Consider, for example, Electron Spin Resonance (ESR).

ESR is used to study materials that have so-called unpaired electrons.

Consider a system of spin- $\frac{1}{2}$ electrons with spin angular momentum $m_s = \pm \frac{1}{2}$

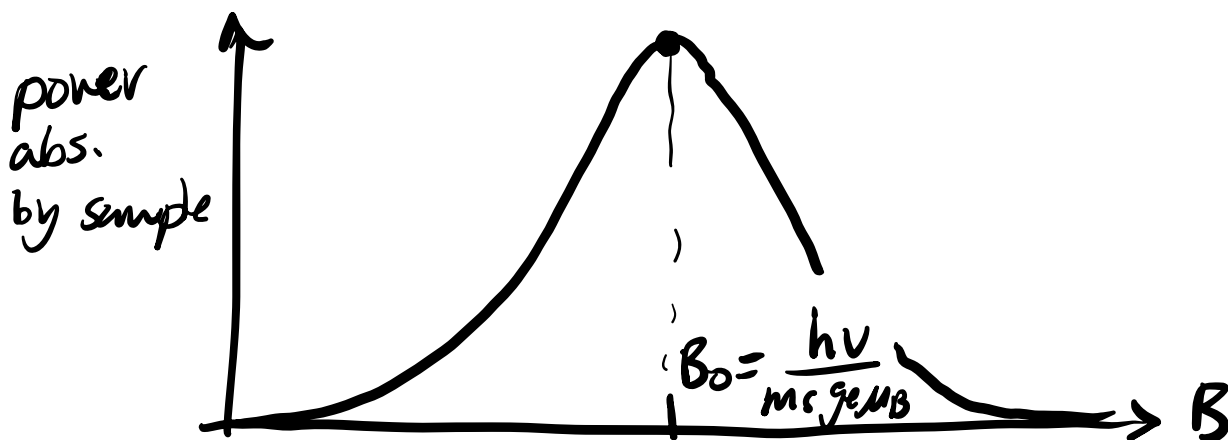
In an external magnetic field applied along z-axis, the energy levels of $m_s = -\frac{1}{2}$

and $m_s = +\frac{1}{2}$ states split.



$E = m_s g_e \mu_B B$ ← magnetic field.
Landé g -factor ↗

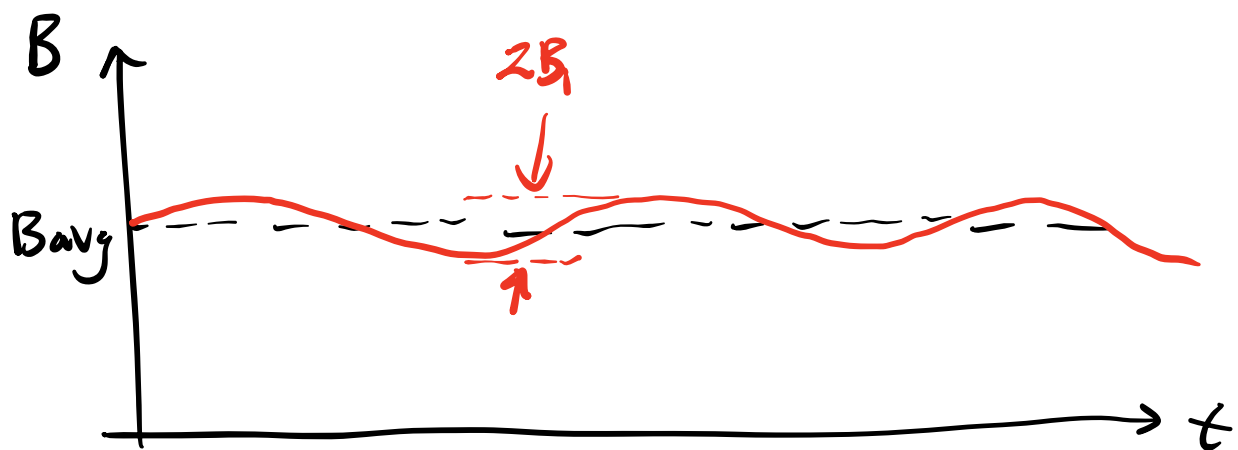
if external field is set to B_0 , then a sample irradiated w/ EM waves w/ energy $h\nu$ can cause transitions between the two states.



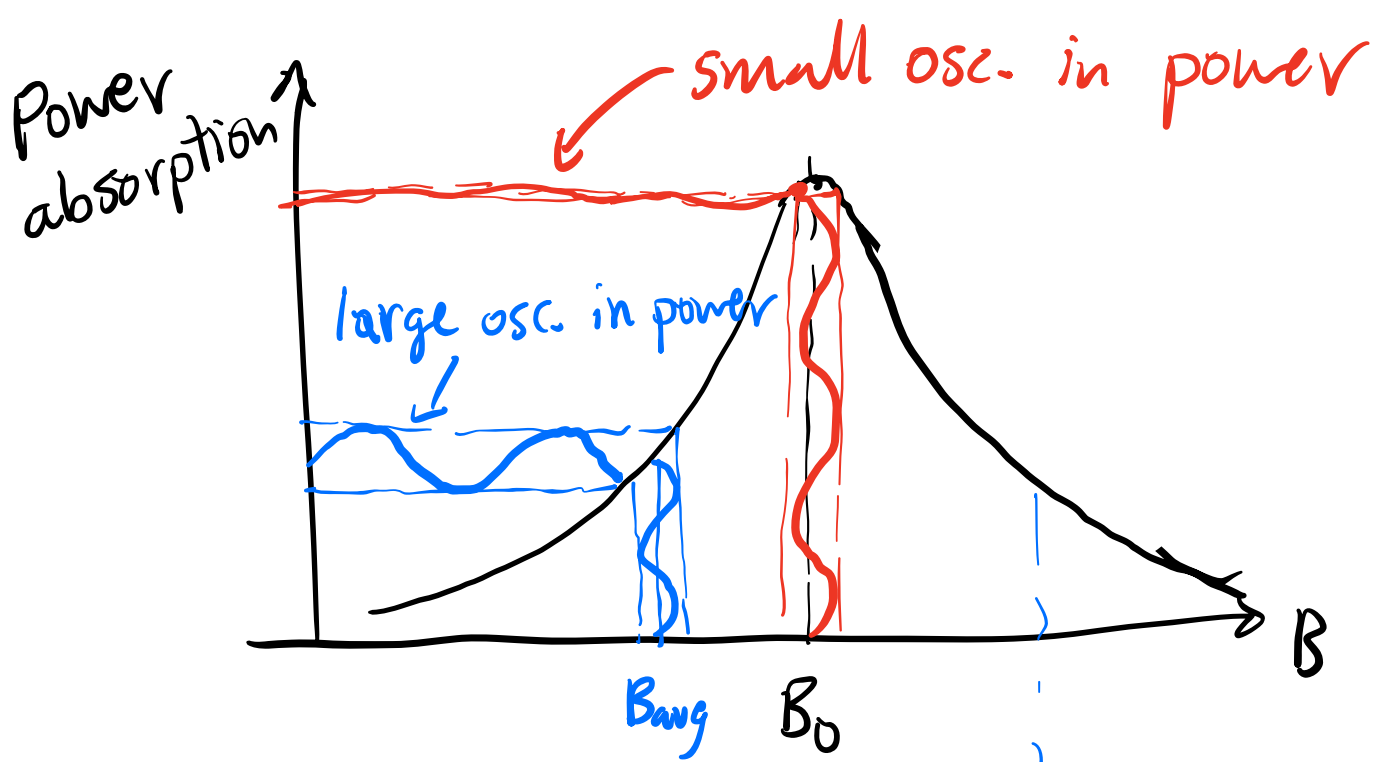
Challenge: The absorption peak can be very weak.

Sol'n: Use lock-in detection.
We need to introduce a sinusoidal signal to the experiment.

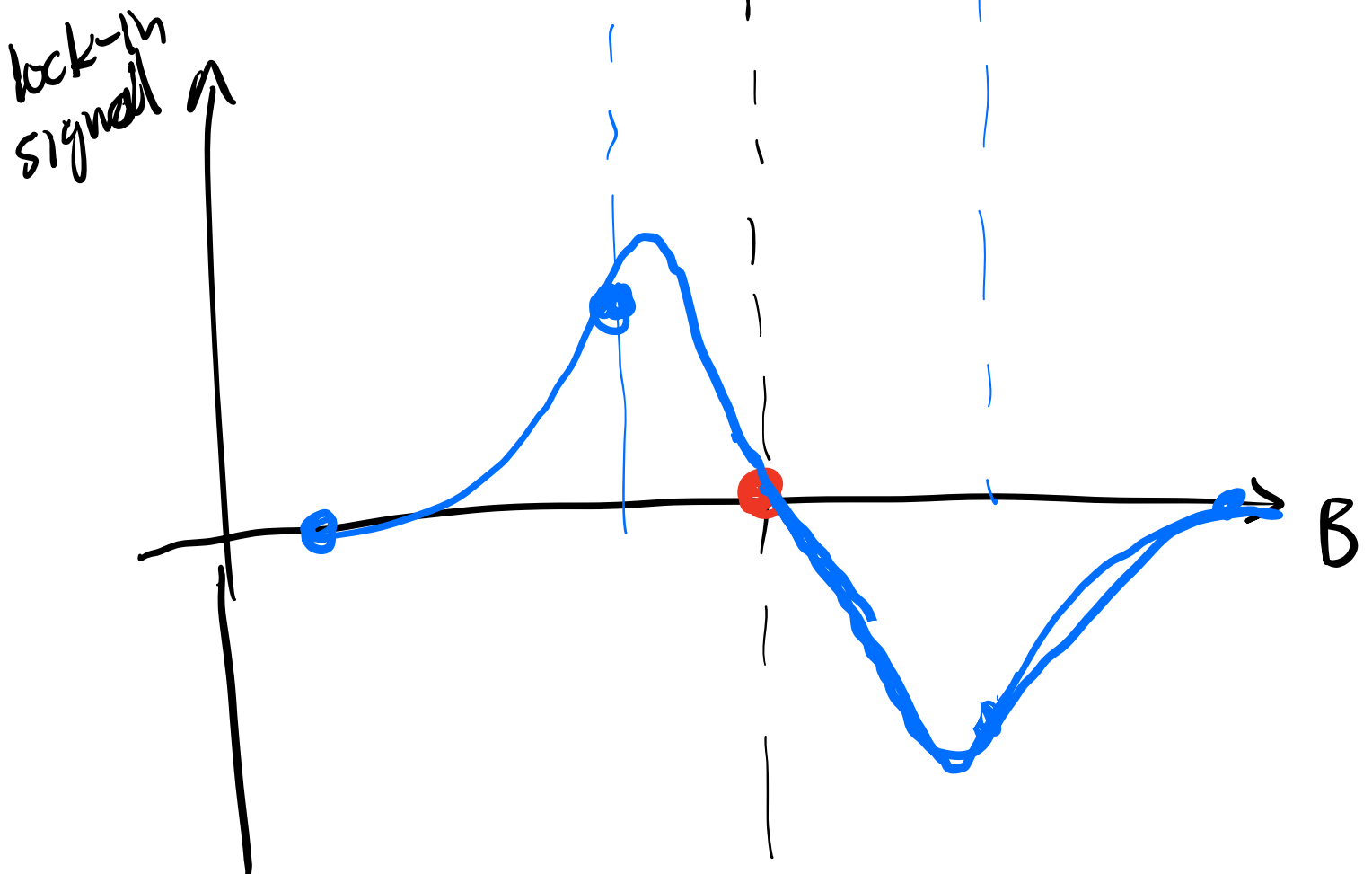
Idea is to add a small AC signal on top of avg. or DC magnetic field.



$$B = B_{avg} + \underbrace{B_1 \sin(\omega t)}_{\text{ref. signal.}}$$



Osc. magnetic field results in an osc. of the power absorption.



Lockin meas. the first derivative of the power absorption signal.